A Note on the Diagonalization of the Discrete Fourier Transform

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Abstract

Following the approach developed by S. Gurevich and R. Hadani, an analytical formula of the canonical basis of the DFT is given for the case $N = p$ where p is a prime number and $p \equiv 1 \pmod 4$.

Index Terms. Discrete Fourier transform, Weil representation, eigenvectors and orthonormal basis.

1 Introduction

The Discrete Fourier transform (DFT) has important applications in communication systems, and can be considered as a family of unitary operators $\{F_N\}$ acting on the Hilbert space $\mathcal{H}_N = \mathbb{C}(\mathbb{Z}_N)$ by the formula

$$F_N[\varphi](j) = \frac{1}{\sqrt{N}} \sum_{i \in \mathbb{Z}_N} e^{\frac{2\pi i}{N} ij} \varphi(i).$$

In the signal processing, the time domain and frequency domain are transformed by the DFT. A canonical basis, in other words, an orthonormal basis of eigenvectors for F_N will simplify the computation of the DFT. The main difficulty to get such a canonical basis is that F_N is an operator of order 4, and it has four distinct eigenvalues $\pm 1, \pm i$ with large multiplicity if $N > 4$. The multiplicity depends on the value of N modulo 4, and was solved in [13], although it was later shown to have been equivalent to a problem solved by Gauss in [5]. Unfortunately, no simple analytical formula for the eigenvectors is known. The research for finding different choices of eigenvectors, selected to satisfy useful properties

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like orthogonality and to have simple forms, has been flourished in the literature [13] [5] [1] [4] [10], just listed a few here.

A novel representation theoretic approach to the diagonalization problem of DFT in the case when $N = p$ is an odd prime number was introduced by S. Gurevich and R. Hadani in [9]. This approach puts to the forefront the Weil representation [17] of the finite symplectic group $Sp = SL_2(\mathbb{F}_p)$ as the fundamental object of underlying harmonic analysis in the finite setting. Specially, a canonical basis Φ_p of eigenvectors of the DFT and the transition matrix Θ_p from the standard basis to Φ_p (*discrete oscillator transform*) for $p \equiv 1 \pmod 4$ were described by an algorithm in [9].

More precisely, a centralizer subgroup G_F of the DFT operator F in $U(\mathcal{H})$ is effectively described by using the Weil representation, then the eigenvectors of G_F are also the eigenvectors of the DFT. The Weil representation in this setting is a unitary representation $\rho : GL_2(\mathbb{F}_p) \to U(\mathcal{H})$ and the DFT is proportional to a single operator $\rho(w)$ where w is the Weyl element. The group G_F is the image under ρ of the centralizer subgroup T_w of w in $GL_2(\mathbb{F}_p)$, where T_w is a maximal algebraic torus (maximal commutative subgroup) in $GL_2(\mathbb{F}_p)$. Restricting the Weil representation to the subgroup T_w yields G_F which commutes with F. A collection of the eigenvectors of the G_F is a canonical basis of the DFT.

The vectors associated to the tori share many nice properties (see [6] [7] [8] for recent applications) and a simple analytical formula for the vectors associated to split tori was given in [16]. Based on the idea of [9] and the approach in [16], an analytical formula of the canonical basis of the DFT for the case of $p \equiv 1 \pmod 4$ is given in this paper. Then the discrete oscillator transform Θ_p introduced in [9] can be obtained in a straightforward manner.

The rest of the paper is organized as follows. In Section 2, we introduce the definitions of the one dimensional finite Heisenberg and Weil representations and the approach studying the eigenvectors of the DFT exhibited in [9]. Then in Section 3, we give an analytical formula of the canonical basis of the DFT, and determine their respective corresponding eigenvalues.

2 Preliminaries

First, we introduce some basic notations which are frequently used in this paper. For a given prime p, let θ and η denote the $(p-1)$th and pth primitive roots of unity in complex field respectively, i.e.,

$$\theta = \exp\left(\frac{2\pi i}{p-1}\right) \quad \text{and} \quad \eta = \exp\left(\frac{2\pi i}{p}\right).$$

We denote $\mathbb{F}_p$ as the finite field with p elements, and $\mathbb{F}_p^*$ as the multiplicative group of $\mathbb{F}_p$ with a generator a. Then for every element $b \in \mathbb{F}_p^*$, there exist i with $0 \leqslant i \leqslant p-2$, such that $b = a^i$. In other words, $i = \log_a b$.

2.1 The Heisenberg Representation

Let (V, ω) be a two-dimensional symplectic vector space over the finite field $\mathbb{F}_p$. For $\forall (t_i, w_i) \in V = \mathbb{F}_p \times \mathbb{F}_p$ $(i = 1, 2)$, the symplectic form ω is given by

$$\omega((t_1, w_1), (t_2, w_2)) = t_1 w_2 - t_2 w_1.$$

Considering V as an Abelian group, it admits a non-trivial central extension called the *Heisenberg group* H. The group H can be presented as $H = V \times F_q$ with the multiplication given by

$$(t_1, w_1, z_1) \cdot (t_2, w_2, z_2) = (t_1 + t_2, w_1 + w_2, z_1 + z_2 + 2^{-1}\omega((t_1, w_1), (t_2, w_2))).$$

It is easy to verify the center of H is $Z = Z(H) = \{(0, 0, z) : z \in \mathbb{F}_p\}$.

For a given non-trivial one dimensional representation ϕ of the center Z, the Heisenberg group H admits a unique irreducible representation of H. We denote $U(\mathcal{H})$ the group of unitary operators on $\mathcal{H}$.

Theorem 1 *(Stone-Von Neuman) Up to isomorphism, there exists a unique irreducible unitary representation $\pi : H \to U(\mathcal{H})$ with central character ϕ, that is, $\pi_{|Z} = \phi \cdot Id_{\mathcal{H}}$.*

The representation π which appears in the above theorem is called the *Heisenberg representation*. In this paper, we take one dimensional representation of Z as $\phi((0, 0, z)) = \eta^z$. Then the unique irreducible unitary representation π corresponding to ϕ has the following formula

$$\pi(t, w, z)[\varphi](i) = \eta^{2^{-1}tw + z + wi}\varphi(i + t) \tag{1}$$

for $\forall \varphi \in \mathbb{C}(\mathbb{F}_p)$, $(t, w, z) \in H$.

2.2 The Weil Representation

The symplectic group $Sp = Sp(V, \omega)$, which is isomorphic to $SL_2(\mathbb{F}_p)$, acts by automorphism of H through its action on the V-coordinate, i.e., for $\forall (t, w, z) \in H$ and a matrix $g = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL_2(\mathbb{F}_p)$, the action g on (t, w, z) is defined as

$$g \cdot (t, w, z) = (at + bw, ct + dw, z). \tag{2}$$

Due to Weil [17], a projective unitary representation $\widetilde{\rho} : Sp \to PGL(\mathcal{H})$ is constructed as follows. Considering the Heisenberg representation $\pi : H \to U(\mathcal{H})$ and $\forall g \in Sp$, a new representation is define as: $\pi^g : H \to U(\mathcal{H})$ by $\pi^g(h) = \pi(g(h))$. Because both π and π^g have the same central character ϕ, they

are isomorphic by Theorem 1. By Schur's Lemma [14], $Hom_H(\pi,\pi^g) \cong \mathbb{C}^*$, so there exist a projective representation $\widetilde{\rho} : Sp \to PGL(\mathcal{H})$. This projective representation $\widetilde{\rho}$ is characterized by the formula:

$$\widetilde{\rho}(g)\pi(h)\widetilde{\rho}(g^{-1}) = \pi(g(h)) \tag{3}$$

for every $g \in Sp$ and $h \in H$. A more delicate statement is that there exists a unique lifting of $\widetilde{\rho}$ into a unitary representation.

Theorem 2 *The projective Weil representation uniquely lifts to a unitary representation*

$$\rho : Sp \to U(\mathcal{H})$$

that satisfies equation (3).

The existence of ρ follows from the fact [3] that any projective representation of $SL_2(\mathbb{F}_p)$ can be lifted to an honest representation, while the uniqueness of ρ follows from the fact [9] that the group $SL_2(\mathbb{F}_p)$ has no non-trivial characters when $p \neq 3$.

Note that $SL_2(\mathbb{F}_p)$ can be generated by $g_a = \begin{pmatrix} a & 0 \\ 0 & a^{-1} \end{pmatrix}$, $g_b = \begin{pmatrix} 1 & 0 \\ b & 1 \end{pmatrix}$, and $w = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$ where $a \in \mathbb{F}_p^*$ and $b \in \mathbb{F}_p$. The formulae of their respective Weil representations for g_a, g_b and w are given in [7] as follows

$$\rho(g_a)[\varphi](i) = \sigma(a)\varphi(a^{-1}i) \tag{4}$$

$$\rho(g_b)[\varphi](i) = \eta^{-2^{-1}bi^2}\varphi(i) \tag{5}$$

$$\rho(w)[\varphi](j) = \frac{1}{\sqrt{p}} \sum_{i \in \mathbb{F}_p} \eta^{ji}\varphi(i) \tag{6}$$

where $\sigma : \mathbb{F}_p^* \to \{\pm 1\}$ is the Legendre character, i.e., $\sigma(a) = a^{\frac{p-1}{2}}$ in $\mathbb{F}_p$.

Obviously, $\rho(w) = F$ is the DFT, and we denote $\rho(g_a) = S_a, \rho(g_b) = N_b$ for convenience. For $\forall g = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL_2(\mathbb{F}_p)$, if $b \neq 0$,

$$g = \begin{pmatrix} a & b \\ c & d \end{pmatrix} = \begin{pmatrix} a & b \\ (ad-1)b^{-1} & d \end{pmatrix} = \begin{pmatrix} b & 0 \\ 0 & b^{-1} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ bd & 1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ ab^{-1} & 1 \end{pmatrix}.$$

Then the Weil representation of g is given by

$$\rho(g) = S_b \circ N_{bd} \circ F \circ N_{ab^{-1}}. \tag{7}$$

If $b = 0$, then

$$g = \begin{pmatrix} a & b \\ c & d \end{pmatrix} = \begin{pmatrix} a & 0 \\ c & a^{-1} \end{pmatrix} = \begin{pmatrix} a & 0 \\ 0 & a^{-1} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ ac & 1 \end{pmatrix}.$$

Hence the Weil representation of g is as follows

$$\rho(g) = S_a \circ N_{ac}. \tag{8}$$

For more details about the Heisenberg and Weil representations, please see [6] [7] [11] [12].

2.3 Centralizer Subgroup of the DFT

A. Maximal Algebraic Tori and T_w

A maximal algebraic *torus* [2] in $SL_2(\mathbb{F}_p)$ is a maximal commutative subgroup which becomes diagonalizable over the original field or quadratic extension of the field. There are two classes of tori in $SL_2(\mathbb{F}_p)$. The first class, called *split tori*, consists of those tori which are diagonalizable over $\mathbb{F}_p$, while the second class, called *non-split tori*, consists of those tori which are not diagonalizable over $\mathbb{F}_p$, but become diagonalizable over the quadratic extension $\mathbb{F}_{p^2}$.

$T_w = \{g : gw = wg, g \in SL_2(\mathbb{F}_p)\}$ is the centralizer group of w in $SL_2(\mathbb{F}_p)$. It is easy to verify

$$T_w = \left\{ \begin{pmatrix} a & -b \\ b & a \end{pmatrix} : a^2 + b^2 = 1, a, b \in \mathbb{F}_p \right\}. \tag{9}$$

If $p \equiv 1 \pmod 4$, then T_w is a split torus and conjugates to the standard diagonal torus

$$A = \left\{ \begin{pmatrix} a & 0 \\ 0 & a^{-1} \end{pmatrix} : a \in \mathbb{F}_p^* \right\}.$$

So T_w is a cyclic subgroup of $SL_2(\mathbb{F}_p)$ with order $p-1$. If $p \equiv 3 \pmod 4$, T_w is a non-split torus which is a cyclic subgroup of $SL_2(\mathbb{F}_p)$ with order $p+1$.

B. Decomposition of Weil representation Associated with T_w

Because T_w is a cyclic group, restricting the Weil representation to T_w: $\rho_{|T_w} : T_w \to U(\mathcal{H})$, we obtains a one dimensional subrepresentation decomposition of $\rho_{|T_w}$ corresponding to an orthogonal decomposition of $\mathcal{H}$(see [14] for basics of group representation theory).

$$\rho_{|T_w} = \bigoplus_{\chi \in \Lambda_{T_w}} \chi \quad \text{and} \quad \mathcal{H} = \bigoplus_{\chi \in \Lambda_{T_w}} \mathcal{H}_\chi \tag{10}$$

where Λ_{T_w} is a collection of all the one dimensional subrepresentation (character) $\chi : T_w \to \mathbb{C}$ in the decomposition of weil representation restricted to T_w.

If $p \equiv 1 \pmod 4$, χ is the character given by $\chi : \mathbb{Z}_{p-1} \to \mathbb{C}$. We have $dim\mathcal{H}_\chi = 1$ unless $\chi = \sigma$ where σ is the Legendre character of T, and $dim\mathcal{H}_\sigma = 2$. If $p \equiv 3 \pmod 4$, χ is the character given by $\chi : \mathbb{Z}_{p+1} \to \mathbb{C}$. There is only one character which does not appear in the decomposition. For the other p characters χ which appear in the decomposition, we have $dim\mathcal{H}_\chi = 1$.

Choosing a generator $t \in T_w$, the character is generated by the eigenvalue $\chi(t)$ of linear operator $\rho(t)$, and the character space $\mathcal{H}_\chi$ naturally corresponds to the eigenspace of $\chi(t)$. Because the eigenvalues of $\rho(t)$ are almost different, it is easier to find a basis of orthogonal eigenvectors of $\rho(t)$ than the DFT. Since $\rho(t)$ commutes with the DFT, the eigenvectors of $\rho(t)$ are also the eigenvectors of the DFT. Thus, we obtain a canonical basis of the DFT.

3 A Canonical Basis of the DFT

If $p \equiv 1 \pmod 4$, T_w is a torus conjugating to the standard diagonal torus $A = \left\{ \begin{pmatrix} a & 0 \\ 0 & a^{-1} \end{pmatrix} : a \in \mathbb{F}_p^* \right\}$, so t, which is a generator of T_w, conjugates to $g_a = \begin{pmatrix} a & 0 \\ 0 & a^{-1} \end{pmatrix}$ where a is a generator of $\mathbb{F}_p^*$, i.e., there exist $s \in SL_2(\mathbb{F}_p)$, such that $t = sg_as^{-1}$ and $\rho(t) = \rho(s)\rho(g_a)\rho(s^{-1})$. Thus, the eigenvectors of $\rho(t)$ can be determined by $\rho(s)$ and the eigenvectors of $\rho(g_a)$. In the following, we first present the results, then their proofs follow.

Lemma 1 *Let* $g_a = \begin{pmatrix} a & 0 \\ 0 & a^{-1} \end{pmatrix}$ *where* a *is a generator of* $\mathbb{F}_p^*$, *then* $\{\psi_x = \{\psi_x(i)\}_{0 \leqslant i < p} : 0 \leqslant x < p\}$ *where*

$$\psi_0(i) = \begin{cases} 1, & for\ i = 0 \\ 0, & for\ i \neq 0 \end{cases} \quad and \quad \psi_x(i) = \begin{cases} 0, & for\ i = 0 \\ \dfrac{1}{\sqrt{p-1}}\theta^{x \log_a i}, & for\ i \neq 0 \end{cases} \quad for \quad 0 < x < p$$

is a orthonormal basis of $\mathcal{H}$ *and a collection of the eigenvectors of* $\rho(g_a)$.

Lemma 2 *Let* $s = \begin{pmatrix} 1 & 2^{-1}a^k \\ a^k & 2^{-1} \end{pmatrix}$ *where* $k = \frac{p-1}{4}$, *then* $t = sg_as^{-1}$ *is a generator of* T_w.

Thus $\Phi_p = \{\varphi_x : \varphi_x = \rho(s)\psi_x, 0 \leqslant x < p\}$ is a canonical basis of $\rho(t) = \rho(sg_as^{-1})$ and the DFT. More explicitly,

Theorem 3 *Let*

$$\varphi_x(i) = \begin{cases} \dfrac{1}{\sqrt{p}}\eta^{2^{-1}a^k i^2}, & for\ x = 0 \\ \dfrac{1}{\sqrt{p(p-1)}}\displaystyle\sum_{j=1}^{p-1}\theta^{x\log_a j}\eta^{a^k(j-i)^2 - 2^{-1}a^k i^2}, & for\ 0 < x < p. \end{cases}$$

Then $\Phi_p = \{\varphi_x = \{\varphi_x(i)\}_{0\leqslant i<p} : 0 \leqslant x < p\}$ *is a orthonormal basis of* $\mathcal{H}$ *and a collection of the eigenvectors of the DFT.*

Theorem 4 $\varphi_x(0 \leqslant x < p)$ *is the eigenvector of the DFT corresponding to the eigenvalue* $(-i)^x$ *where* $i = \sqrt{-1}$*, i.e.,*

$$F\varphi_x = (-i)^x\varphi_x.$$

Now we prove the above lemmas and theorems. Considering $\{\delta_i : i \in \mathbb{F}_p\}$ which is the orthonormal basis of Hilbert space $\mathcal{H} = \mathbb{C}(\mathbb{F}_p)$, where δ_i is defined as $\delta_i(j) = \delta_{ij}$ for $\forall i, j \in \mathbb{F}_p$, every vector $\varphi = \{\varphi(i)\}$ can be written as the form $\varphi = \sum_{i\in\mathbb{F}_p}\varphi(i)\delta_i$. Recall that $SL_2(\mathbb{F}_p)$ can be generated by $g_a = \begin{pmatrix} a & 0 \\ 0 & a^{-1} \end{pmatrix}$, $g_b = \begin{pmatrix} 1 & 0 \\ b & 1 \end{pmatrix}$ and $w = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$ where $a \in \mathbb{F}_p^*$ and $b \in \mathbb{F}_p$. Then their respective Weil representations (4),(5), and (6) of g_a, g_b, and w can be rewritten as follows

$$\rho(g_a)\delta_i = S_a\delta_i = \sigma(a)\delta_{ai} \tag{11}$$

$$\rho(g_b)\delta_i = N_b\delta_i = \eta^{-2^{-1}bi^2}\delta_i \tag{12}$$

$$\rho(w)\delta_j = F\delta_j = \frac{1}{\sqrt{p}}\sum_{i\in\mathbb{F}_p}\eta^{ji}\delta_i. \tag{13}$$

Proof of Lemma 1. From (11), we have

$$\rho(g_a)\delta_i = \sigma(a)\delta_{ai} = -\delta_{ai}.$$

Let $V_1 = V(\delta_1), V_2 = V(\delta_2, \delta_3, \cdots, \delta_{p-1})$, then it is obvious that $\mathcal{H} = V_1 \bigoplus V_2$, $< V_1, V_2 >= 0$, and $\rho(g_a)(V_i) = V_i$ for $i = 1, 2$. It is easy to see that $\rho(g_a)_{|V_1} = -Id$, so δ_0 is a eigenvector of $\rho(g_a)$ corresponding to the eigenvalue -1. The eigenfunction of $\rho(g_a)_{|V_2}$ is $(x^{p-1} - 1)$, so the eigenvalues of $\rho(g_a)_{|V_2}$ are $\theta^0, \theta^1, \theta^2 \cdots\cdots \theta^{p-2}$ which are different. We assert that $\sum_{i=1}^{p-1}\theta^{(\frac{p-1}{2}-j)\log_a i}\delta_i$ is the

eigenvector associated to the eigenvalue $\theta^j (0 \leqslant j \leqslant p-2)$, and it can be verified as follows

$$\begin{aligned}
\rho(g_a)(\sum_{i=1}^{p-1} \theta^{(\frac{p-1}{2}-j)\log_a i}\delta_i) &= -\sum_{i=1}^{p-1} \theta^{(\frac{p-1}{2}-j)\log_a i}\delta_{ai} \\
&= -\sum_{i=1}^{p-1} \theta^{(\frac{p-1}{2}-j)\log_a (a^{-1}i)}\delta_i \\
&= \theta^{\frac{p-1}{2}} \sum_{i=1}^{p-1} \theta^{(\frac{p-1}{2}-j)(\log_a i - 1)}\delta_i \\
&= \theta^{\frac{p-1}{2}}\theta^{j-\frac{p-1}{2}} \sum_{i=1}^{p-1} \theta^{(\frac{p-1}{2}-j)\log_a i}\delta_i \\
&= \theta^{j} \sum_{i=1}^{p-1} \theta^{(\frac{p-1}{2}-j)\log_a i}\delta_i.
\end{aligned}$$

Let $x = \frac{q-1}{2} - j$. By normalizing the eigenvectors, we complete the proof. □

Proof of Lemma 2. Note that

$$\begin{aligned}
t &= sg_a s^{-1} \\
&= \begin{pmatrix} 1 & 2^{-1}a^k \\ a^k & 2^{-1} \end{pmatrix} \begin{pmatrix} a & 0 \\ 0 & a^{-1} \end{pmatrix} \begin{pmatrix} 2^{-1} & -2^{-1}a^k \\ -a^k & 1 \end{pmatrix} \\
&= \begin{pmatrix} 2^{-1}(a-a^{-1}) & -2^{-1}a^k(a-a^{-1}) \\ 2^{-1}a^k(a-a^{-1}) & 2^{-1}(a-a^{-1}) \end{pmatrix} \in T_w.
\end{aligned}$$

On the other hand, t conjugates to g_a, so the order of t is $p-1$. Thus, t is a generator of T_w.

Proof of Theorem 3. Since $t = sg_a s^{-1}$, $\Phi_p = \{\varphi_x = \rho(s)\psi_x : 0 \leqslant x < p\}$ is a collection of the orthogonal eigenvectors of $\rho(t)$ where φ_x and s are presented in Lemmas 1 and 2 respectively.

From (12), s has the following decomposition

$$s = \begin{pmatrix} 1 & 2^{-1}a^k \\ a^k & 2^{-1} \end{pmatrix} = \begin{pmatrix} 2^{-1}a^k & 0 \\ 0 & 2a^{3k} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 4^{-1}a^k & 1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 2a^{3k} & 1 \end{pmatrix}.$$

Then applying (11),(12), and (13), for $1 \leqslant x \leqslant p-1$, we have

$$\begin{aligned}
\varphi_x = \rho(s)\psi_x &= S_{2^{-1}a^k} \circ N_{4^{-1}a^k} \circ F \circ N_{2a^{3k}}\Big(\frac{1}{\sqrt{p-1}} \sum_{j=1}^{p-1} \theta^{x\cdot \log_a j} \delta_j\Big) \\
&= S_{2^{-1}a^k} \circ N_{4^{-1}a^k} \circ F\Big(\frac{1}{\sqrt{p-1}} \sum_{j=1}^{p-1} \theta^{x\cdot \log_a j} \eta^{a^k j^2} \delta_j\Big) \\
&= S_{2^{-1}a^k} \circ N_{4^{-1}a^k}\Big(\frac{1}{\sqrt{p(p-1)}} \sum_{i=0}^{p} \sum_{j=1}^{p-1} \theta^{x\cdot \log_a j} \eta^{a^k j^2 + ij} \delta_i\Big) \\
&= S_{2^{-1}a^k}\Big(\frac{1}{\sqrt{p(p-1)}} \sum_{i=0}^{p} \sum_{j=1}^{p-1} \theta^{x\cdot \log_a j} \eta^{a^k j^2 + ij - 8^{-1}a^k i^2} \delta_i\Big) \\
&= \frac{\sigma(2^{-1}a^k)}{\sqrt{p(p-1)}} \sum_{i=0}^{p} \sum_{j=1}^{p-1} \theta^{x\cdot \log_a j} \eta^{a^k j^2 + ij - 8^{-1}a^k i^2} \delta_{2^{-1}a^k i}) \quad \text{(substitute } i \text{ by } 2^{-1}a^k i) \\
&= \frac{\sigma(2^{-1}a^k)}{\sqrt{p(p-1)}} \sum_{i=0}^{p} \sum_{j=1}^{p-1} \theta^{x\cdot \log_a j} \eta^{a^k j^2 + 2a^{-k} ij + 2^{-1}a^k i^2} \delta_i \\
&= \frac{\sigma(2^{-1}a^k)}{\sqrt{p(p-1)}} \sum_{i=0}^{p} \sum_{j=1}^{p-1} \theta^{x\cdot \log_a j} \eta^{a^k (j-i)^2 - 2^{-1}a^k i^2} \delta_i.
\end{aligned}$$

For $x = 0$, we have

$$\varphi_0 = \rho(s)\phi_0 = \frac{\sigma(2^{-1}a^k)}{\sqrt{p}} \sum_{i=0}^{p} \eta^{2^{-1}a^k i^2} \delta_i.$$

Because $\{\psi_x : 0 \leqslant x < p\}$ is a orthonormal basis and $\rho(s)$ is a unitary matrix, $\{\varphi_x : 0 \leqslant x < p\}$ is also a orthonormal basis of $\mathcal{H}$. Since $\rho(s)F = F\rho(s)$, $\varphi_x (0 \leqslant x < p)$ are not only the eigenvectors of $\rho(s)$, but also the eigenvectors of the DFT. Note that $\sigma(2^{-1}a^k)$ is a constant, which completes the proof. □

Proof of Theorem 4. It can be verified as follows, for $x = 0$, we have

$$\begin{aligned}
F[\varphi_0](t) &= \frac{1}{p}\sum_{t=0}^{p-1}\eta^{2^{-1}a^k i^2+it} \\
&= \frac{1}{p}\eta^{2^{-1}a^k t^2}\sum_{t=0}^{p-1}\eta^{2^{-1}a^k i^2+it-2^{-1}a^k t^2} \\
&= \frac{1}{p}\eta^{2^{-1}a^k t^2}\sum_{t=0}^{p-1}(\eta^{2^{-1}a^k})^{(i-a^k t)^2} \\
&= \frac{1}{p}\eta^{2^{-1}a^k t^2}\sum_{t=0}^{p-1}(\eta^{2^{-1}a^k})^{t^2} \quad \text{(substitute } i-a^k t \text{ by } t\text{)} \\
&= \frac{1}{\sqrt{p}}\eta^{2^{-1}a^k t^2} \quad \text{(Gauss sum)} \\
&= \varphi_0(t).
\end{aligned}$$

For $x \neq 0$, we have

$$\begin{aligned}
F[\varphi_x](t) &= \frac{1}{p\sqrt{p-1}}\sum_{t=0}^{p-1}\sum_{j=1}^{p-1}\theta^{x\log_a j}\eta^{a^k(j-i)^2-2^{-1}a^k i^2+it} \\
&= \frac{1}{p\sqrt{p-1}}\sum_{j=1}^{p-1}\theta^{x\log_a j}\sum_{t=0}^{p-1}\eta^{a^k(j-i)^2-2^{-1}a^k i^2+it} \\
&= \frac{1}{p\sqrt{p-1}}\sum_{j=1}^{p-1}\theta^{x\log_a a^{-k}j}\sum_{t=0}^{p-1}\eta^{a^k(a^{-k}j-i)^2-2^{-1}a^k i^2+it} \quad \text{(substitute } j \text{ by } a^{-k}j\text{)} \\
&= \frac{\theta^{-kx}}{p\sqrt{p-1}}\sum_{j=1}^{p-1}\theta^{x\log_a j}\sum_{t=0}^{p-1}\eta^{a^{-k}j^2-2ji+2^{-1}a^k i^2+it} \\
&= \frac{(-i)^x}{p\sqrt{p-1}}\sum_{j=1}^{p-1}\theta^{x\log_a j}\eta^{a^{-k}j^2-2^{-1}a^{-k}(t-2j)^2}\sum_{t=0}^{p-1}\eta^{2^{-1}a^k i^2+i(t-2j)+2^{-1}a^{-k}(t-2j)^2} \\
&= \frac{(-i)^x}{p\sqrt{p-1}}\sum_{j=1}^{p-1}\theta^{x\log_a j}\eta^{a^k(j-t)^2-2^{-1}a^k t^2}\sum_{t=0}^{p-1}\eta^{(2a^k)^{-1}(a^k i+t-2j)^2} \\
&= \frac{(-i)^x}{\sqrt{p(p-1)}}\sum_{j=1}^{p-1}\theta^{x\log_a j}\eta^{a^k(j-t)^2-2^{-1}a^k t^2} \quad \text{(Gauss sum)} \\
&= (-i)^x\varphi_x(t).
\end{aligned}$$

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